

One Step Forward, Two Steps Back: Vector Potential in SPMHD



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Smoothed Particle Magnetohydrodynamics

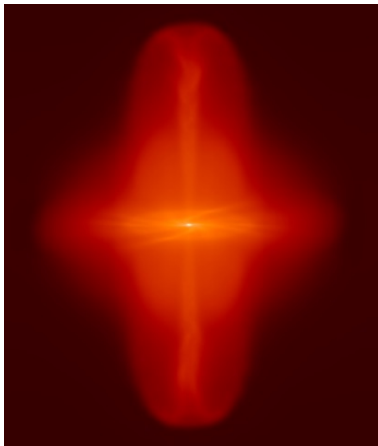
- SPMHD simulates magnetohydrodynamics (magnetized fluids).
- Magnetic field is coupled to the hydrodynamic equations.

$$\frac{dv^i}{dt} = \frac{1}{\rho} \frac{\partial S^{ij}}{\partial x^j} \quad S^{ij} = -\delta^{ij} \left(P + \frac{B^2}{2\mu_0} \right) + \frac{B^i B^j}{\mu_0}$$

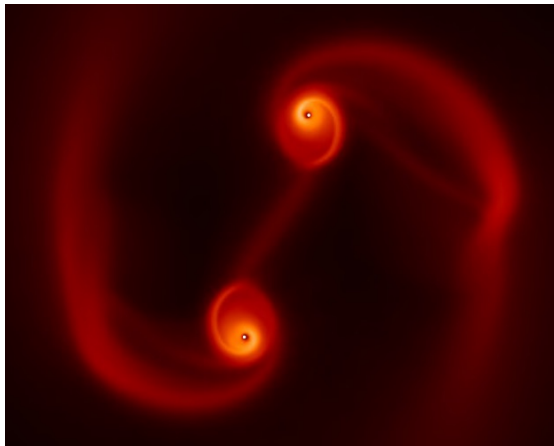
$$\frac{d\mathbf{B}}{dt} = -(\mathbf{B} \cdot \nabla) \mathbf{v} + \mathbf{B} (\nabla \cdot \mathbf{v})$$

$$\nabla \cdot \mathbf{B} = 0$$

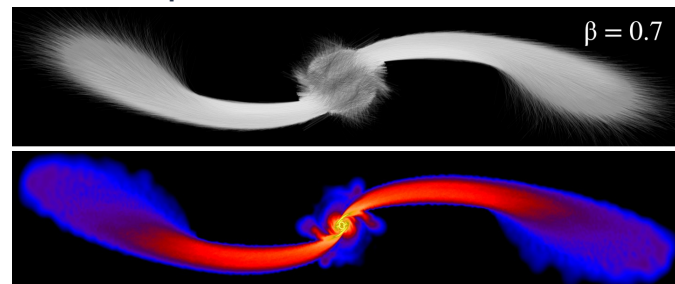
Star Formation



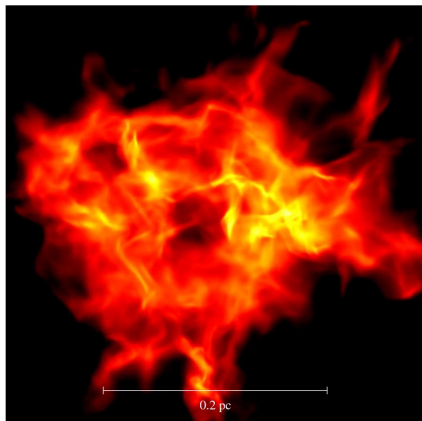
Binary Star Formation



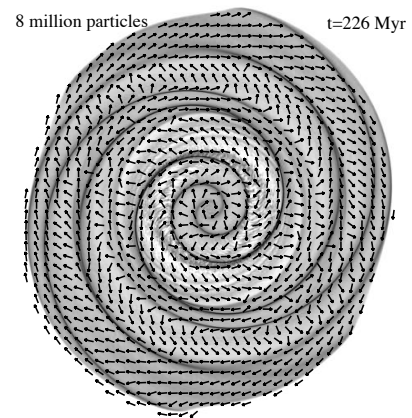
Tidal Disruption Events



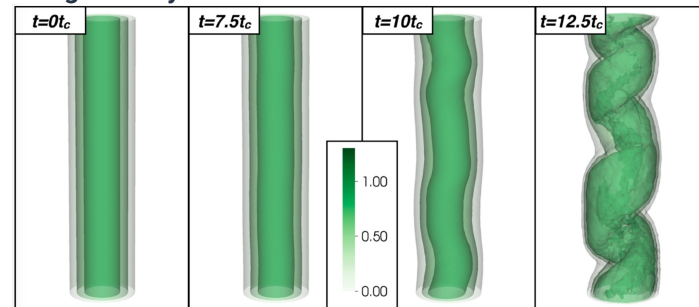
Star Cluster Formation



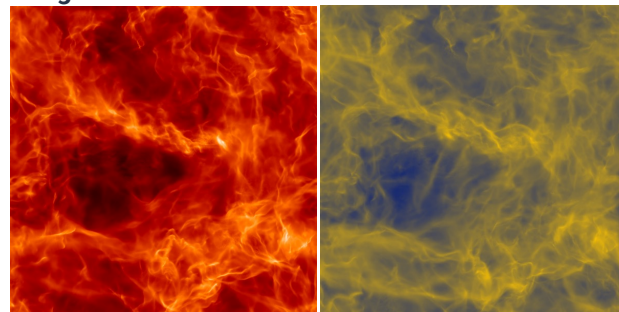
Milky Way



Magnetically Confined Plasmas



Magnetized Turbulence



Price, **Tricco** & Bate (2012); Wurster, Price, Bate (2016); Liptai et al (2017); Dobbs, Price, Pettitt, Bate, **Tricco** (2016); Bonnerot et al (2017); Vela Vela et al (2019); **Tricco** Price & Federrath (2016)

Divergence-Free Condition

- The primary challenge in simulating magnetized fluids with SPH is the divergence-free condition.

$$\nabla \cdot \mathbf{B} = 0$$

- For the past decade, SPMHD simulations have used **constrained hyperbolic divergence cleaning** (Tricco, Price 2012; Tricco, Price, Bate 2016).
- Constrained cleaning typically provides a **10x reduction in divergence error**, and keeps average **error around the ~1% level**.

Goal: Create an exactly divergence-free SPMHD

Vector Potential

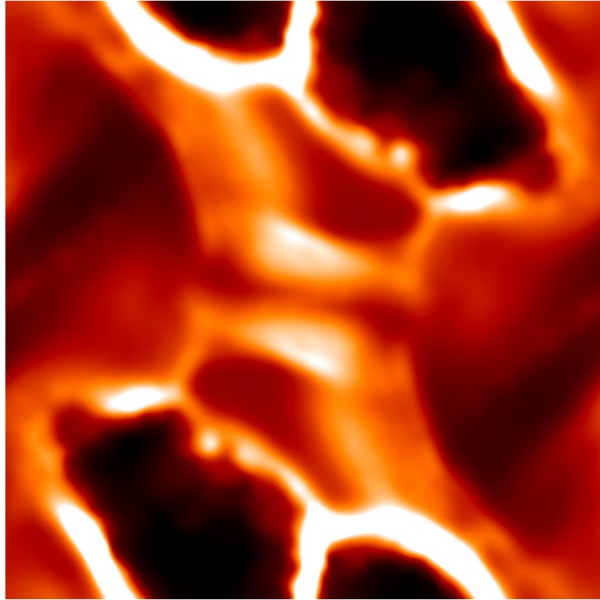
- Formulating the magnetic field in terms of the **Vector Potential** is a logical choice.

$$\mathbf{B} = \nabla \times \mathbf{A}$$

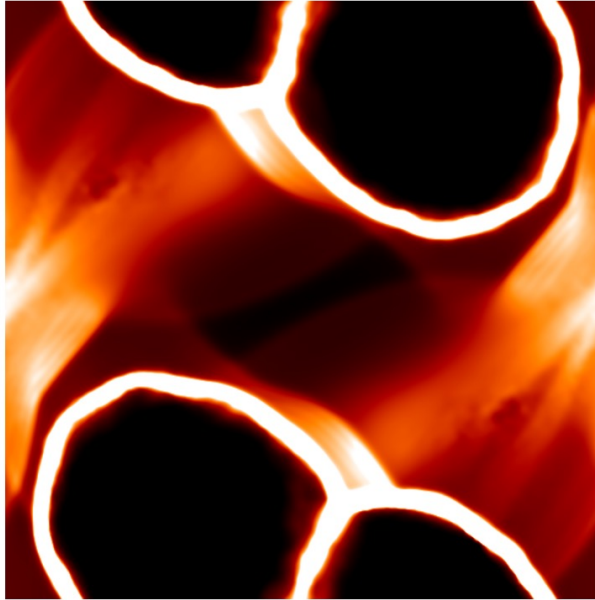
- The divergence of the curl is zero.
- Guaranteed to create a divergence-free magnetic field by construction.
- Price (2010) investigated the vector potential, finding **severe numerical issues**.

Orszag-Tang Vortex with Vector Potential

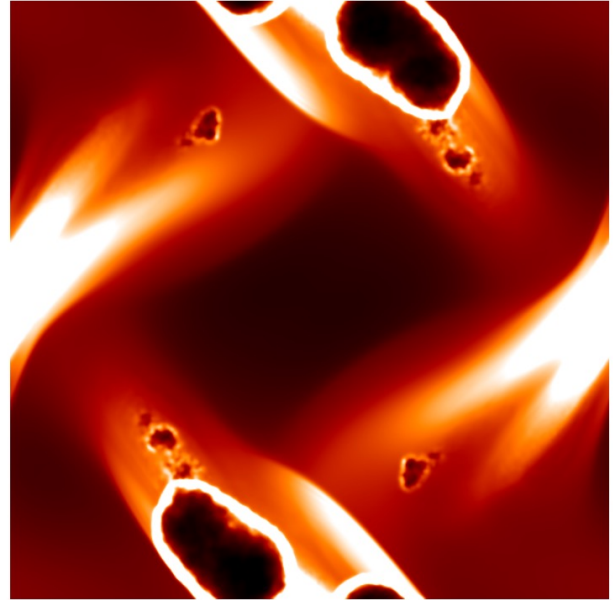
$n_x = 64$ particles, $t = 0.5$



$n_x = 128$ particles, $t = 0.3$



$n_x = 256$ particles, $t = 0.1975$



Price (2010) Formulation

Start with the induction equation and use $\mathbf{B} = \nabla \times \mathbf{A}$.

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{v} \times \mathbf{B}) \quad \longrightarrow \quad \frac{\partial \mathbf{A}}{\partial t} = \mathbf{v} \times (\nabla \times \mathbf{A}) + \nabla \phi$$

Choose the gauge $\phi = -\mathbf{v} \cdot \mathbf{A}$ (constant of integration) yields the Galilean invariant discretised equations:

$$\frac{dA_a^i}{dt} = \frac{A_a^j}{\Omega_a \rho_a} \sum_b m_b \left(v_a^j - v_b^j \right) \frac{\partial W_{ab}(h_a)}{\partial x_a^i}$$

A New Hope?

Our new approach is to express the vector potential in **integral form**.

$$\begin{aligned}\frac{\partial \mathbf{A}}{\partial t} = \mathbf{v} \times (\nabla \times \mathbf{A}) \quad \longrightarrow \quad \frac{d}{dt} \int_V \mathbf{A} dV &= \int_v \mathbf{v} \times \mathbf{B} dV + \int_{\partial V} \mathbf{A} \mathbf{v} \cdot d\mathbf{S}, \\ &= \int_v \mathbf{v} \times \mathbf{B} dV + \int_V \nabla_j (\mathbf{A} \mathbf{v}^j) dV.\end{aligned}$$

Which discretized into SPMHD using $\phi = -\mathbf{v} \cdot \mathbf{A}$ yields

$$\begin{aligned}\frac{dA_a^i}{dt} &= \frac{A_a^j}{\Omega_a \rho_a} \sum_b m_b (v_a^j - v_b^j) \frac{\partial W_{ab}(h_a)}{\partial x_a^i} \\ &\quad + \frac{1}{\Omega_a \rho_a} \sum_b m_b (A_a^i - A_b^i) \left[(v_a^j - v_b^j) \frac{\partial W_{ab}(h_a)}{\partial x_a^j} \right]\end{aligned}$$

Comparison

Price (2010) formulation

$$\frac{dA_a^i}{dt} = \frac{A_a^j}{\Omega_a \rho_a} \sum_b m_b \left(v_a^j - v_b^j \right) \frac{\partial W_{ab}(h_a)}{\partial x_a^i}$$

New integral formulation

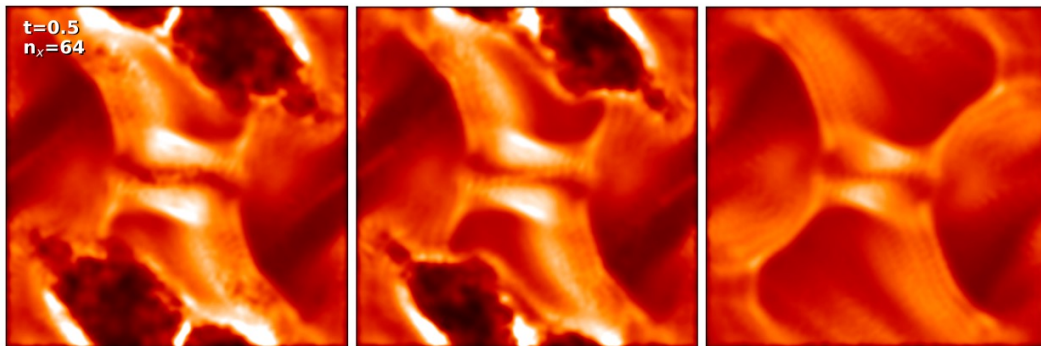
$$\begin{aligned} \frac{dA_a^i}{dt} = & \frac{A_a^j}{\Omega_a \rho_a} \sum_b m_b \left(v_a^j - v_b^j \right) \frac{\partial W_{ab}(h_a)}{\partial x_a^i} \\ & + \frac{1}{\Omega_a \rho_a} \sum_b m_b \left(A_a^i - A_b^i \right) \left[\left(v_a^j - v_b^j \right) \frac{\partial W_{ab}(h_a)}{\partial x_a^j} \right] \end{aligned}$$

VOLUME INTEGRAL

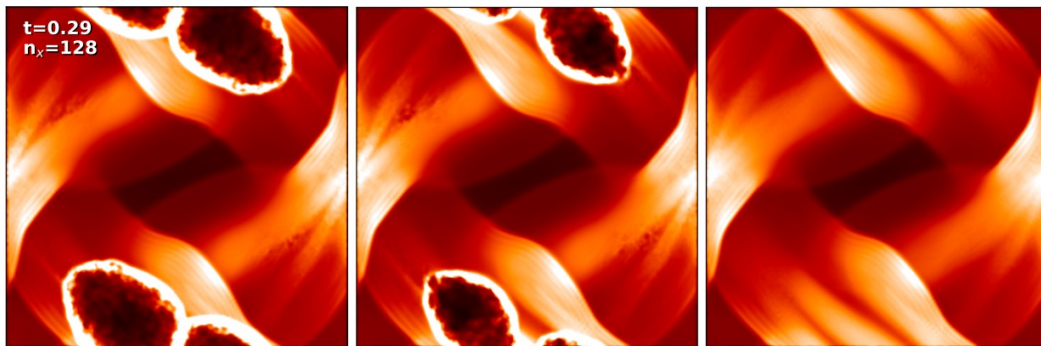
PRICE (2010)

DIRECTLY EVOLVING **B**

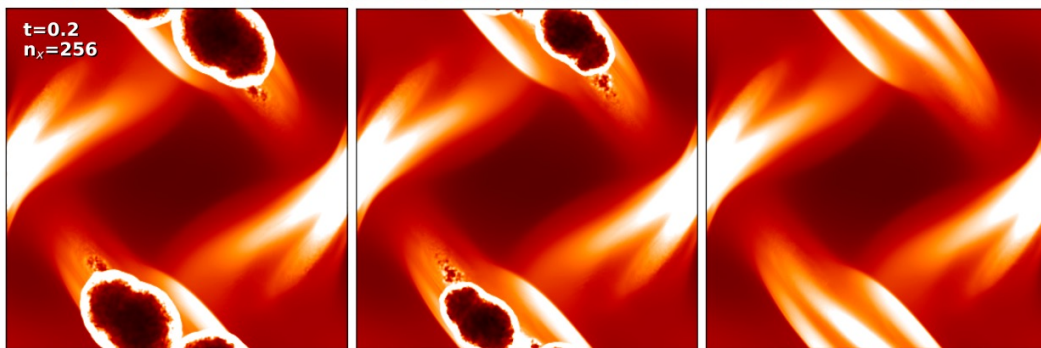
$n_x = 64$ particles, $t = 0.5$



$n_x = 128$ particles, $t = 0.29$



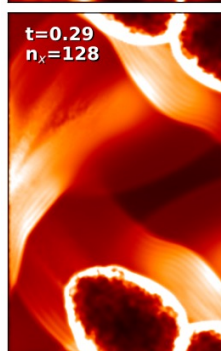
$n_x = 256$ particles, $t = 0.2$



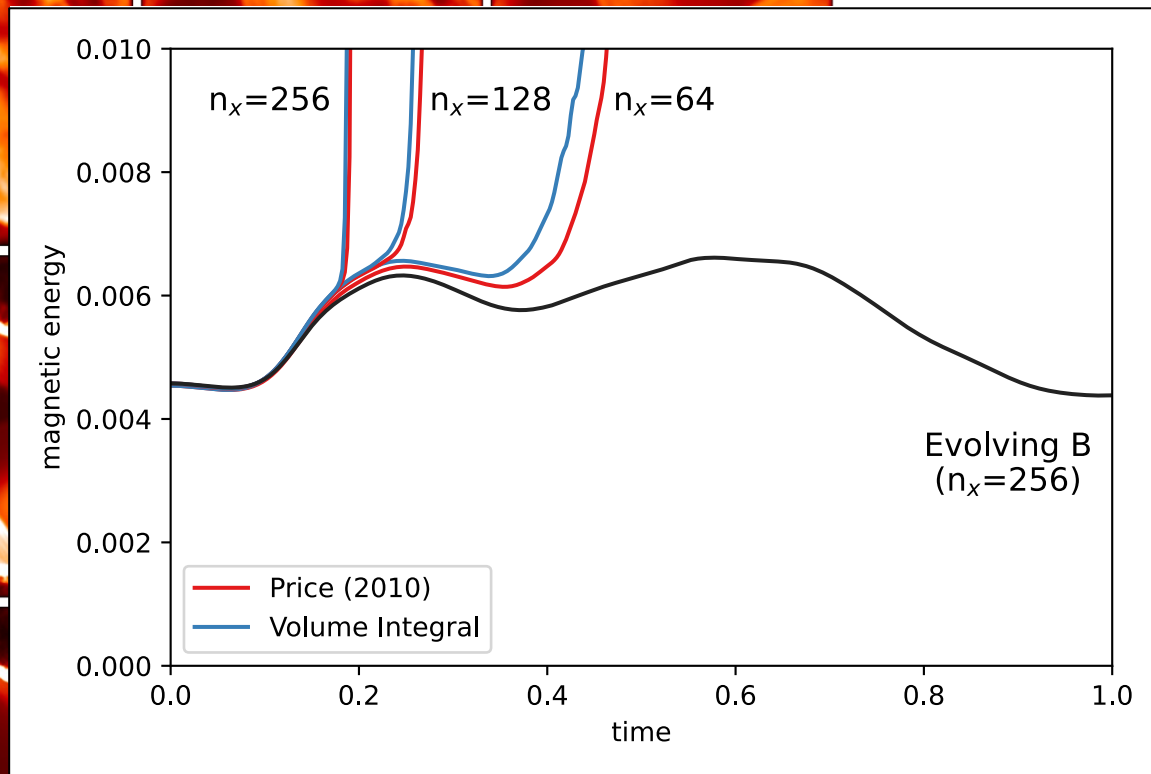
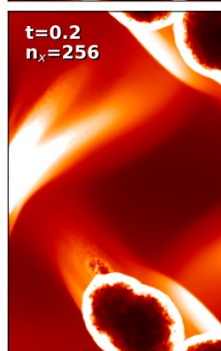
$n_x = 64$ particles, $t = 0.5$



$n_x = 128$ particles, $t = 0.29$



$n_x = 256$ particles, $t = 0.2$



Vector Potential in SPMHD

- The integral vector potential formulation is still highly numerically unstable.
- **Why?**
 - One issue is that the magnetic field is evolved using **A**, but the equations of motion use **B**.
 - This mismatch of variables does not conserve energy, leading to exponential energy growth and numerical instability.

A Fully Conservative Formulation

- Enforcing energy conservation can be achieved by solving for the equations of motion from the Lagrangian (specified in terms of **A**, not **B**).

$$\begin{aligned}
 \frac{dv_a^i}{dt} = & - \sum_b m_b \left[\frac{P_a}{\rho_a^2} + \frac{P_b}{\rho_b^2} \right] \frac{\partial W_{ab}}{\partial x_a^i} && \longrightarrow \text{Pressure gradient} \\
 & + \frac{3}{2\mu_0} \sum_b m_b \left[\frac{B_a^2}{\rho_a^2} + \frac{B_b^2}{\rho_b^2} \right] \frac{\partial W_{ab}}{\partial x_a^i} \\
 & - \frac{1}{\mu_0} \epsilon_{jkl} \sum_b m_b (A_a^k - A_b^k) \left[\frac{B_a^j}{\rho_a^2} + \frac{B_b^j}{\rho_b^2} \right] \frac{\partial^2 W_{ab}}{\partial x_a^i \partial x_a^l} && \left. \begin{array}{l} \\ \\ \end{array} \right\} \text{Equivalent to Price (2010) Formulation} \\
 & - \sum_b m_b \left[\frac{A_a^i}{\rho_a^2} J_a^k + \frac{A_b^i}{\rho_b^2} J_b^k \right] \frac{\partial W_{ab}}{\partial x_a^k} \\
 & - \sum_b m_b (A_a^k - A_b^k) \left[\frac{J_a^k}{\rho_a^2} + \frac{J_b^k}{\rho_b^2} \right] \frac{\partial W_{ab}}{\partial x_a^i} && \longrightarrow \text{Additional term from integral approach}
 \end{aligned}$$

A Fully Conservative Formulation

- Most problematic is that the equations of motion **do not represent the MHD equations of motion in the continuum limit!**
- Continuum equations derive from $\mathbf{J} \times \mathbf{B}$ where $\mathbf{J} = \nabla \times \mathbf{B} / \mu_0$

$$\frac{dv^i}{dt} = -\frac{1}{\rho} \frac{\partial P}{\partial x^i} + \frac{1}{\rho} \left[J^j \frac{\partial A^j}{\partial x^i} - J^j \frac{\partial A^i}{\partial x^j} \right]$$

- The integral approach equations of motion are equivalent to (note the spurious factor of 3)

$$\frac{dv^i}{dt} = -\frac{1}{\rho} \frac{\partial P}{\partial x^i} + \frac{1}{\rho} \left[3J^j \frac{\partial A^j}{\partial x^i} - J^j \frac{\partial A^i}{\partial x^j} \right]$$

Summary

- An **integral approach** was developed to model the magnetic vector potential.
- Key results:
 - A correction term to the Price (2010) vector potential evolution equation is introduced. Results are **not improved**.
 - The conservative equations of motion were derived, which includes a correction term to the Price (2010) approach.
 - These equations of motion **do not represent** the MHD equations.
- A truly divergence-free SPMHD solution continues to remain elusive.